

# Testing Supersymmetry with Lepton Flavor Violating $\tau$ and $\mu$ decays

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# Motivation

- ★ Strong evidence of lepton flavor changing neutrino oscillations in neutrino data
  - Neutrinos have non-zero masses
  - Lepton number violating processes can occur (Leptogenesis...)
  - Lepton flavor violating processes can occur ( $\mu \rightarrow e\gamma$ ...)
  
- ★ Lepton Flavor Violating processes: Window to look for new physics
  - There are not LFV processes within SM ( $m_\nu = 0$ )
  
- ★ We assume Majorana  $\nu_R$  responsible for  $m_\nu \neq 0$
  
- ★ Seesaw mechanism for neutrino mass generation
  
- ★ Large LFV rates in SUSY-seesaw models due to flavor mixing in soft slepton masses
  - many studies on  $\mu \rightarrow e\gamma$ ,  $\tau$  rare decays,  $Z \rightarrow \tau\bar{\mu}$ ...
  
- ★ We study LFV decays of  $\tau$  and  $\mu$  within the CMSSM-seesaw:
  - $\tau^- \rightarrow \mu^- \mu^- \mu^+$ ,  $\tau^- \rightarrow e^- e^- e^+$ ,  $\mu^- \rightarrow e^- e^- e^+$
  - Full one-loop computation
  - Parallel study of  $l_j \rightarrow l_i \gamma$ ,  $i \neq j$ , to insure compatibility with data

## LFV in the MSSM-seesaw

- We assume universal soft-SUSY-breaking masses at large energies  $M_X \gg m_M$ ,

$$m_M = \text{Majorana } \nu_R \text{ mass; } M_X = 2 \times 10^{16} \text{ GeV}$$

$$(m_{\tilde{L}})_{ij}^2 = M_0^2 \delta_{ij}, (m_{\tilde{E}})_{ij}^2 = M_0^2 \delta_{ij}, (m_{\tilde{M}})_{ij}^2 = M_0^2 \delta_{ij},$$

$$M_1 = M_2 = M_{1/2}, M_{H_1} = M_{H_2} = M_0$$

- The RGE-running down to  $m_M$  generates flavor non-diagonal soft masses due to  $Y_\nu$ ,

$$(\Delta m_{\tilde{L}}^2)_{ij} = -\frac{1}{8\pi^2}(3M_0^2 + A_0^2)(Y_\nu^* L Y_\nu^T)_{ij}$$

$$(\Delta A_l)_{ij} = -\frac{3}{16\pi^2} A_0 Y_{l_i} (Y_\nu^* L Y_\nu^T)_{ij}$$

$$(\Delta m_{\tilde{E}}^2)_{ij} = 0 ; L_{kl} \equiv \log \left( \frac{M_X}{m_{M_k}} \right) \delta_{kl}$$

- We assume three  $\nu_R$  and use the parametrization of A. Casas and A. Ibarra (Nucl.Phys.B 618(2001)171)

Flavor diagonal basis for  $l$  and  $\nu_R$ ; Flavor non-diagonal  $Y_\nu$ ,

$$Y_{l_i} = \frac{m_{l_i}}{v_1}, (Y_\nu)_{ij} = \frac{(m_D)_{ij}}{v_2}, v_{(1,2)} = v(\cos \beta, \sin \beta)$$

$$m_D^T = i m_N^{diag 1/2} R m_\nu^{diag 1/2} U_{MNS}^+ ; R^T R = 1$$

$R$  complex matrix;  $m_N^{diag}$ ,  $m_\nu^{diag}$  physical masses

# Seesaw and physical parameters setup

★ We consider two scenarios, compatible with data:

- quasi-degenerate light and degen. heavy neutrinos:

$$m_{\nu_1} = 0.2 \text{ eV}, m_{\nu_2} = m_{\nu_1} + \frac{\Delta m_{sol}^2}{2m_{\nu_1}}, m_{\nu_3} = m_{\nu_1} + \frac{\Delta m_{atm}^2}{2m_{\nu_1}},$$

$$m_{N_1} = m_{N_2} = m_{N_3} = m_N$$

- hierarchical light and hierarchical heavy neutrinos:

$$m_{\nu_1} \simeq 0 \text{ eV}, m_{\nu_2} = \sqrt{\Delta m_{sol}^2}, m_{\nu_3} = \sqrt{\Delta m_{atm}^2},$$

$$m_{N_1} \leq m_{N_2} < m_{N_3}$$

★ We take input values for  $\Delta m^2$  and  $U_{MNS}$  :

$$\sqrt{\Delta m_{sol}^2} = 0.008 \text{ eV}, \sqrt{\Delta m_{atm}^2} = 0.05 \text{ eV}; \delta = \alpha = \beta = 0$$

$$\theta_{12} = \theta_{sol} = 30^\circ, \theta_{23} = \theta_{atm} = 45^\circ, \theta_{13} = 0^\circ,$$

★ We compute the LFV  $\tau$  and  $\mu$  decay rates for various choices of  $m_{N_{1,2,3}}$  and  $R$

$$R = \begin{pmatrix} c_2 c_3 & -c_1 s_3 - s_1 s_2 c_3 & s_1 s_3 - c_1 s_2 c_3 \\ c_2 s_3 & c_1 c_3 - s_1 s_2 s_3 & -s_1 c_3 - c_1 s_2 s_3 \\ s_2 & s_1 c_2 & c_1 c_2 \end{pmatrix}$$

$$c_i = \cos \theta_i; s_i = \sin \theta_i, \theta_{1,2,3} = \text{complex angles}$$

Parameters in  $R$  constrained by perturbativity of  $Y_\nu$

Favourable values for Baryogenesis  $m_{N_{1,2,3}} > 10^8 \text{ GeV}$

# MSSM parameters setup

★ We work in the physical basis:

Involved SUSY particles in the one-loop diagrams:  $\chi^-, \chi^0, \tilde{l}, \tilde{\nu}$

● Charge slepton sector:

Once leptons and neutrinos are rotated to the physical basis, the slepton mass matrices are NOT diagonal:

**lepton-slepton misalignment**

$$M_{\tilde{l}}^2 = \begin{pmatrix} M_{LL}^{ee2} & M_{LR}^{ee2} & M_{LL}^{e\mu2} & M_{LR}^{e\mu2} & M_{LL}^{e\tau2} & M_{LR}^{e\tau2} \\ M_{RL}^{ee2} & M_{RR}^{ee2} & M_{RL}^{e\mu2} & M_{RR}^{e\mu2} & M_{RL}^{e\tau2} & M_{RR}^{e\tau2} \\ M_{LL}^{\mu e2} & M_{LR}^{\mu e2} & M_{LL}^{\mu\mu2} & M_{LR}^{\mu\mu2} & M_{LL}^{\mu\tau2} & M_{LR}^{\mu\tau2} \\ M_{LR}^{\mu e2} & M_{RR}^{\mu e2} & M_{RL}^{\mu\mu2} & M_{RR}^{\mu\mu2} & M_{RL}^{\mu\tau2} & M_{RR}^{\mu\tau2} \\ M_{LL}^{\tau e2} & M_{LR}^{\tau e2} & M_{LL}^{\tau\mu2} & M_{LR}^{\tau\mu2} & M_{LL}^{\tau\tau2} & M_{LR}^{\tau\tau2} \\ M_{RL}^{\tau e2} & M_{RR}^{\tau e2} & M_{RL}^{\tau\mu2} & M_{RR}^{\tau\mu2} & M_{RL}^{\tau\tau2} & M_{RR}^{\tau\tau2} \end{pmatrix}$$

referred to  $(\tilde{e}_L, \tilde{e}_R, \tilde{\mu}_L, \tilde{\mu}_R, \tilde{\tau}_L, \tilde{\tau}_R)$

$$M_{LL}^{ij2} = m_{\tilde{L},ij}^2 + \frac{1}{2}v_1^2 (Y_l^\dagger Y_l)_{ij} + m_Z^2 \cos 2\beta \left( -\frac{1}{2} + \sin^2 \theta_W \right) \delta_{ij}$$

$$M_{RR}^{ij2} = m_{\tilde{E},ij}^2 + \frac{1}{2}v_1^2 (Y_l^\dagger Y_l)_{ij} - m_Z^2 \cos 2\beta \sin^2 \theta_W \delta_{ij}$$

$$M_{LR}^{ij2} = \frac{1}{\sqrt{2}} \left( v_1 (A_l^{ij})^* - \mu Y_l^{ij} \right)$$

$$M_{RL}^{ij2} = (M_{LR}^{ij2})^*, \quad i, j = e, \mu, \tau$$

● Sneutrino sector: ( $\tilde{\nu}_R$  decouple)

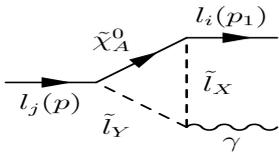
$$M_{\tilde{\nu}}^2 = \begin{pmatrix} m_{\tilde{L},e}^2 + \frac{1}{2}m_Z^2 \cos 2\beta & m_{\tilde{L},e\mu}^2 & m_{\tilde{L},e\tau}^2 \\ m_{\tilde{L},\mu e}^2 & m_{\tilde{L},\mu}^2 + \frac{1}{2}m_Z^2 \cos 2\beta & m_{\tilde{L},\mu\tau}^2 \\ m_{\tilde{L},\tau e}^2 & m_{\tilde{L},\tau\mu}^2 & m_{\tilde{L},\tau}^2 + \frac{1}{2}m_Z^2 \cos 2\beta \end{pmatrix}$$

referred to  $(\tilde{\nu}_{eL}, \tilde{\nu}_{\mu L}, \tilde{\nu}_{\tau L})$

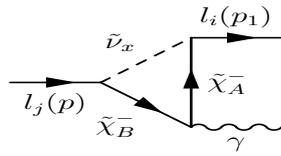
★ Soft parameters from RGE-running down to EW scale (SPheno programme)

★  $\mu$  from EW breaking condition

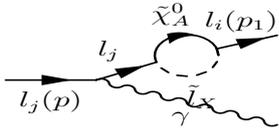
# $\gamma$ -penguin diagrams



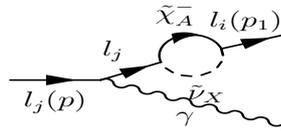
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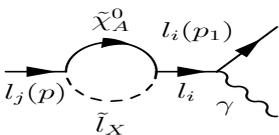
(G2)



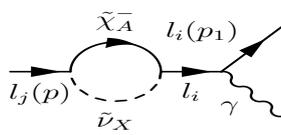
(G3)



(G4)

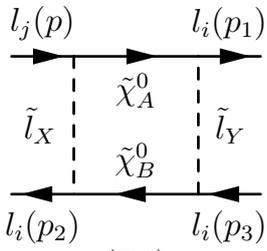


(G5)

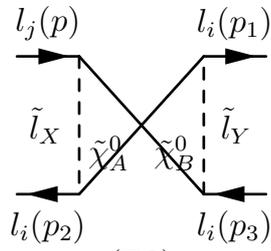


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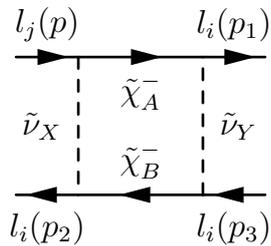
# Box-type diagrams



(B1)

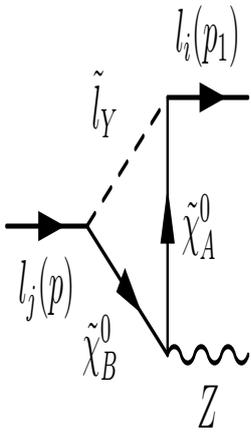


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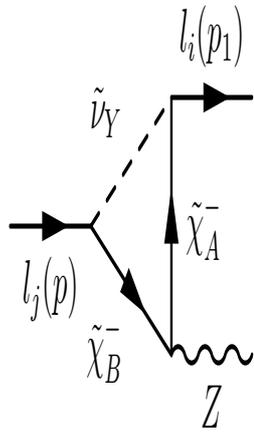


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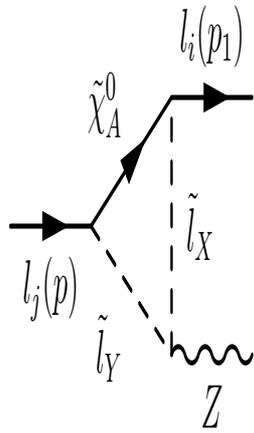
# Z-penguin diagrams



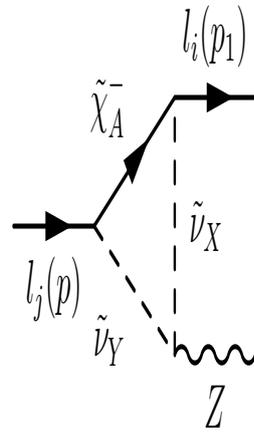
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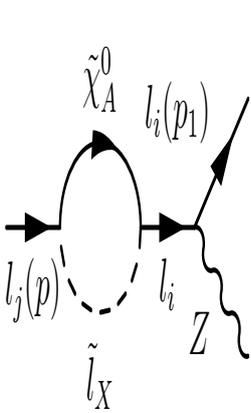
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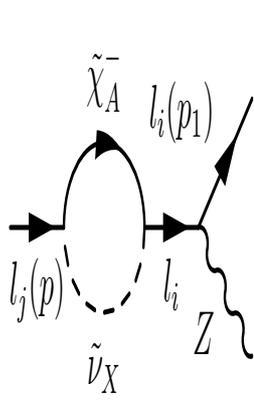
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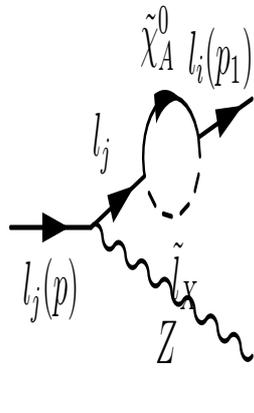
(Z4)



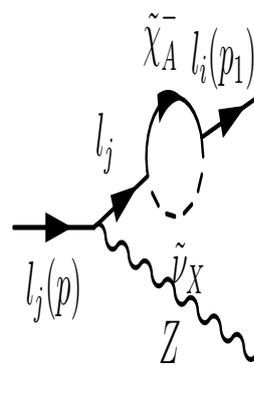
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(Z6)

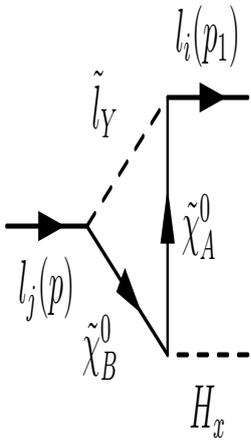


(Z7)

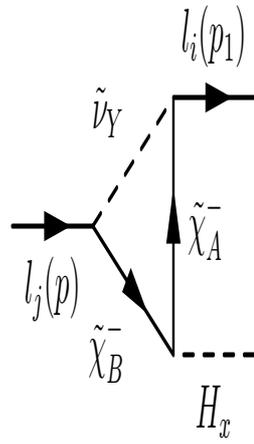


(Z8)

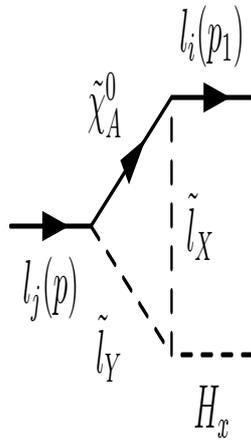
# H-penguin diagrams



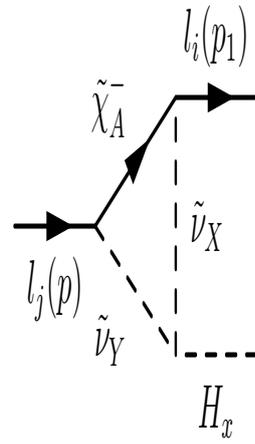
(Z1)



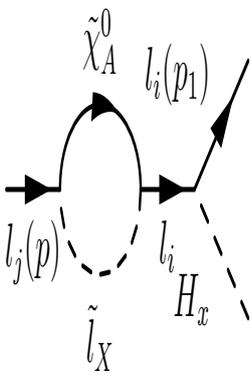
(Z2)



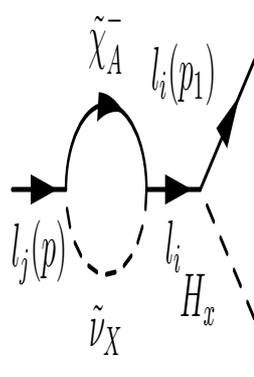
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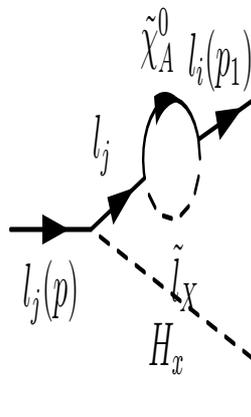
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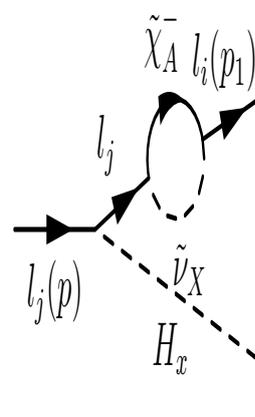
(Z5)



(Z6)



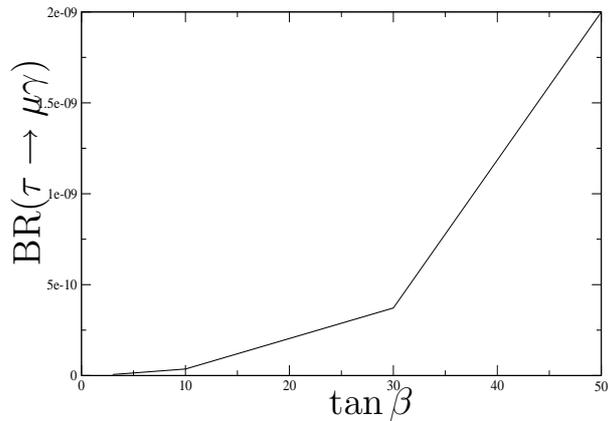
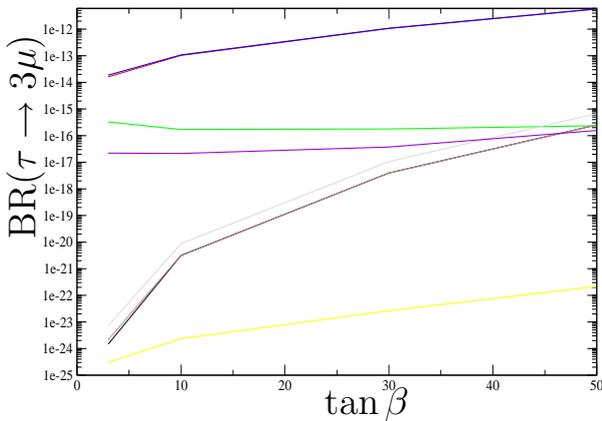
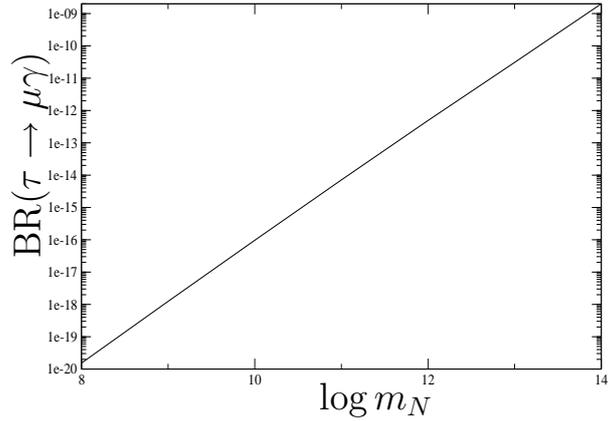
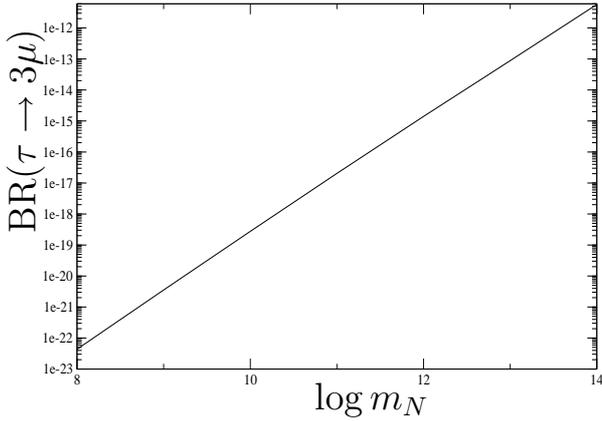
(Z7)



(Z8)

# Branching ratios dependence with $m_N$ and $\tan \beta$ degenerate $m_{N_i}$ and real $R$

Explored values:  $10^8 \leq m_N(\text{GeV}) \leq 10^{14}$   
 $(M_0, M_{1/2}, A_0) = (400, 300, 0)$  GeV



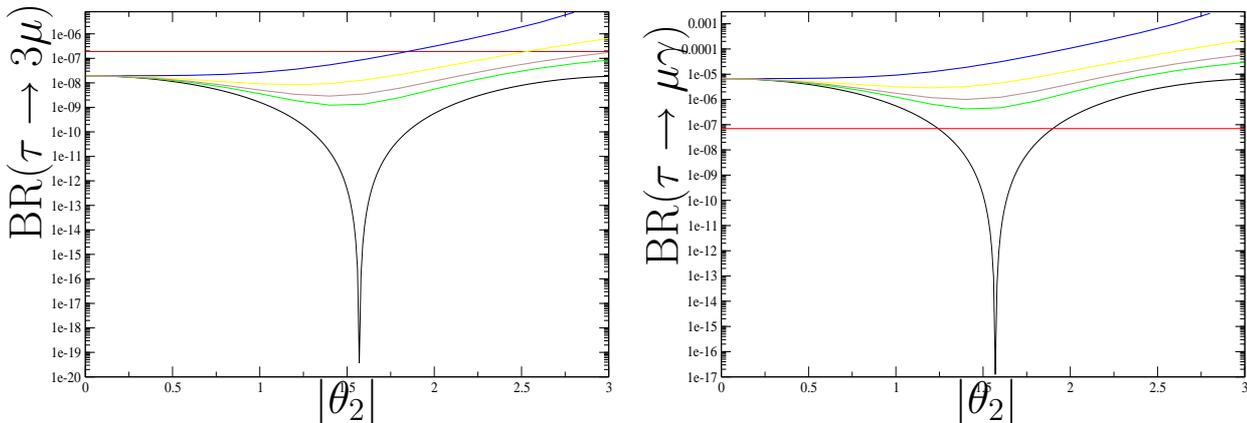
- ★ All rates grow with  $m_N(\text{GeV})$  as expected
- ★ Total rates grow as  $(\tan \beta)^2$
- ★  $BR(\tau^- \rightarrow \mu^- \mu^- \mu^+) < 10^{-11}$ ,  $\tan \beta < 50$
- ★  $BR(\tau \rightarrow \mu\gamma)$  below exp. limits:  $7 \times 10^{-8}$  for  $m_N < 10^{14}$  GeV
- ★ Smaller rates for  $\tau \rightarrow 3e$ ,  $\mu \rightarrow 3e$ ,  $\tau \rightarrow e\gamma$ ,  $\mu \rightarrow e\gamma$

# Branching ratios dependence with $R$ hierarchical neutrinos

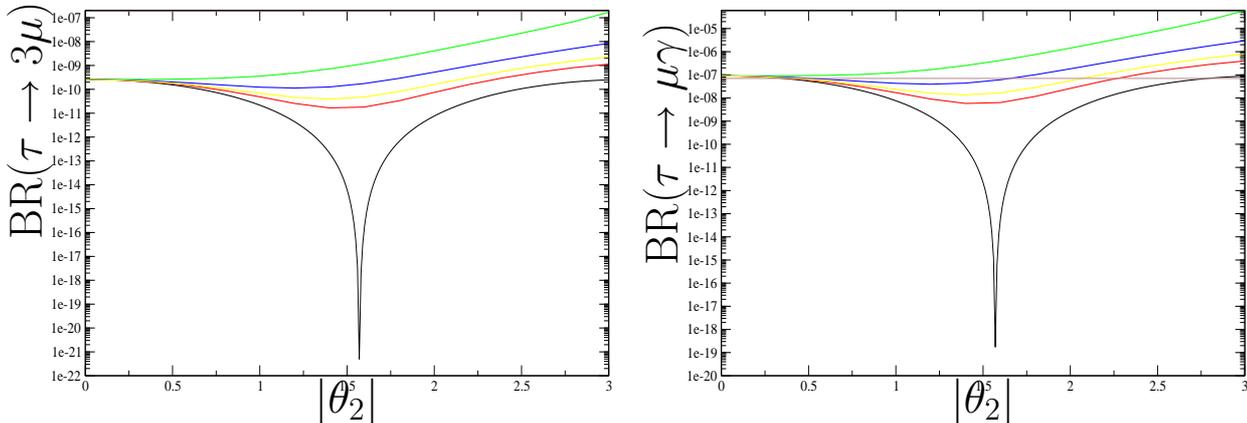
$$(m_{N_1}, m_{N_2}, m_{N_3}) = (10^8, 2 \times 10^8, 10^{14}) \text{ GeV}$$

★ **Complex  $\theta_2$**  ( $\arg(\theta_2) = 0, \pi/4, \pi/6, \pi/8, \pi/10, \theta_1 = \theta_3 = 0$ )

$$\tan \beta = 50, (M_0, M_{1/2}, A_0) = (200, 100, 0) \text{ GeV}$$



$$\tan \beta = 50, (M_0, M_{1/2}, A_0) = (400, 300, 0) \text{ GeV}$$



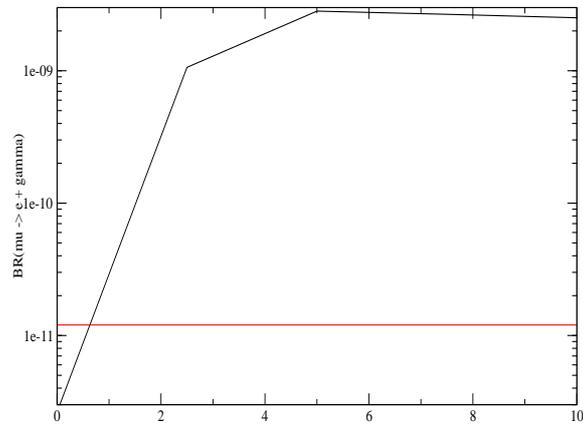
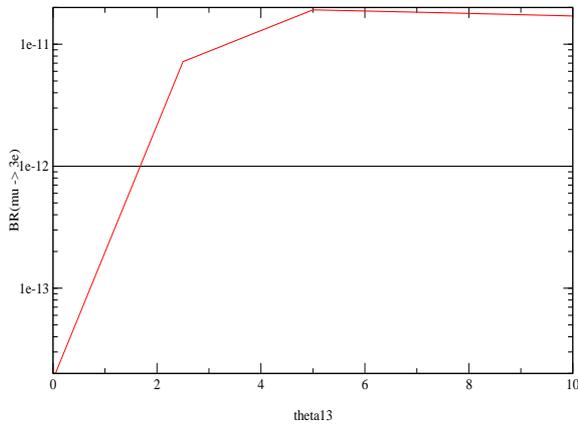
**Hierarchical:** Much larger rates than for degenerate case.  
**Light SUSY spectra** severely constrained

# Branching ratios dependence with $\theta_{13}$ hierarchical neutrinos

$$(m_{N_1}, m_{N_2}, m_{N_3}) = (10^8, 2 \times 10^8, 10^{14}) \text{ GeV}$$

$$\theta_1 = \theta_2 = \theta_3 = 0$$

$$\tan \beta = 50, (M_0, M_{1/2}, A_0) = (400, 300, 0) \text{ GeV}$$



$\theta_{13} > 2^\circ$  are unallowed by LFV data

## Conclusions:

- ★ LFV  $\tau$  and  $\mu$  decays with ratios  $BR(\tau^- \rightarrow \mu^- \mu^- \mu^+) \sim 10^{-5}$  are found in the MSSM-seesaw for hierarchical neutrinos with the largest heavy mass at  $10^{14}$  GeV and large  $\tan \beta = 50$
- ★ SUSY parameters significantly restricted
- ★ Deserves further studies